

TD6 Conduction thermique à symétrie cylindrique et effet Joule

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Convection Examen Janvier 2011

Thermodynamique statistique

① A)

1) $P_J = \sigma_U \times V =$ puissance dissipée par effet Joule

$$\sigma_U = \frac{P_J}{V} = \frac{R I^2}{V} = \frac{\rho_c L}{S_1} \times \frac{I^2}{S_1 L} = \frac{\rho_c I^2}{\pi^2 b_1^4}$$

$$\sigma_U = 15502 \text{ W m}^{-3}$$

2) $0 \leq r \leq b_1$ Loi de Fourier $\vec{J}_U = -\lambda_c \frac{dT}{dr} \vec{e}_r$
(J_U indépendant de ϕ et z)

$$3) \frac{\partial(\rho u)}{\partial t} = -\operatorname{div} \vec{J}_U + \nabla u = 0 \quad (\text{régime stationnaire})$$

$$-\operatorname{div}(-\lambda_c \vec{\text{grad}} T) + \nabla u = 0 ; \quad \lambda_c \Delta T + \nabla u = 0$$

$$\Delta T = \left\{ \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) \right\} = -\frac{\sigma_U}{\lambda_c}$$

$$\frac{\partial}{\partial r} \left(r \frac{\partial T}{\partial r} \right) = -\frac{\sigma_U}{\lambda_c} \Rightarrow r \frac{\partial T}{\partial r} = -\frac{\sigma_U}{\lambda_c} \frac{r^2}{2} + A$$

$$\frac{dT}{dr} = -\frac{\sigma_U}{2\lambda_c} r + \frac{A}{r} \quad \text{avec } \begin{cases} A=0 \\ \text{sinon } \frac{dT}{dr} \rightarrow \infty \text{ en } r=0 \end{cases}$$

$$T(r) = -\frac{\sigma_U}{4\lambda_c} r^2 + T(r=0) = -\frac{\rho_c I^2}{4\pi^2 b_1^4 \lambda_c} r^2 + T_0$$

$$T(r=b_1) = T_c = -\frac{\rho_c I^2}{4\pi^2 b_1^4 \lambda_c} + T_0 = T_0$$

$\frac{4\pi^2 b_1^4 \lambda_c}{110^{-3}}$ négligeable

$\frac{dT}{dr}$ du conducteur négligeable car conductivité thermique faible ($390 \text{ W m}^{-1} \text{ K}^{-1}$)

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$$4) \quad I'_U(r) = \iint_{S_2} \vec{J}'_U \cdot \vec{n} \, dS = J'_U 2\pi r L$$

avec $b_1 \leq r \leq b_2$

$$J'_U(r) = \frac{I'_U(r)}{2\pi r L} = -\lambda_i \frac{dT}{dr}$$

$$I_U(r) = I'_U(r) = P_J = \frac{\rho_c I^2}{\pi b_1^2} L$$

~~($0 \leq r \leq b_1$) ($b_1 \leq r \leq b_2$)~~

$$\frac{dT}{dr} = -\frac{\rho_c I^2}{\pi b_1^2} \frac{L}{\lambda_i 2\pi r L} = -\frac{\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \frac{1}{r}$$

$$T(r) = \frac{-\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \ln r + B$$

$$T(r=b_2) = T_a = -\frac{\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \ln b_2 + B$$

$$B = T_a + \frac{\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \ln b_2$$

$$\boxed{T(r) = T_a + \frac{\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \ln \left(\frac{b_2}{r} \right)}$$

$$5) \quad T(r=b_1) = T_a + \frac{\rho_c I^2}{2\pi^2 b_1^2 \lambda_i} \ln \left(\frac{b_2}{b_1} \right) = T_c$$

$$\lambda_i = \frac{\rho_c I^2}{2\pi^2 b_1^2 (T_c - T_a)} \ln \left(\frac{b_2}{b_1} \right) = 0,1 \text{ W m}^{-1}\text{K}^{-1}$$

⑤ 6) ~~$P_J = J_U \times V = J_U \times \pi b_1^2 \times L$~~

~~↓ transmet par convection de la surface latérale du fil à la surface latérale du conducteur extérieur~~